

$$3. \det: O(n) \rightarrow \mathbb{Z}_2 \quad AA^T = \mathbf{1} \Rightarrow \det = \pm 1$$

$$M \mapsto \det(M)$$

$$\ker(\det) = SO(n)$$

$$1 \rightarrow SO(n) \rightarrow O(n) \rightarrow \mathbb{Z}_2 \rightarrow 1$$

$$4. \det: U(n) \rightarrow U(1)$$

$$U \mapsto \det U = \lambda \quad \ker? \quad SU(n)$$

$$1 \rightarrow SU(n) \rightarrow U(n) \rightarrow U(1) \rightarrow 1$$

$$5. \pi: SU(2) \rightarrow SO(3)$$

$$U \vec{x} \cdot \vec{\sigma} U^\dagger := (\pi(U) \vec{x}) \cdot \vec{\sigma}$$

$$U \in \ker \pi \quad U \vec{x} \cdot \vec{\sigma} U^\dagger = \vec{x} \cdot \vec{\sigma} \quad U = \lambda \mathbf{1}$$

$$\lambda = \pm 1$$

$$\pi(U) = \pi(-U)$$

$$\Rightarrow \ker \pi \cong \mathbb{Z}_2$$

$$SU(2)/\mathbb{Z}_2 \cong SO(3)$$

$$1 \rightarrow \mathbb{Z}_2 \rightarrow SU(2) \rightarrow SO(3) \rightarrow 1$$

↑
abelian, and $\mathbb{Z}_2 \subset Z(SU(2))$

Definition (central extension)

$$1 \rightarrow A \xrightarrow{i} E \xrightarrow{\pi} G \rightarrow 1$$

1. A is abelian.

$$2. A \subset Z(E) \quad i(a)b = b i(a) \quad (\forall a \in A, \forall b \in E)$$

$$1 \rightarrow N \rightarrow G^{\text{Quantum}} \rightarrow G^{\text{Classical}} \rightarrow 1$$

Motivation for such extensions. in QM $\begin{cases} (1.1) \\ (1.5.1 \& 2) \end{cases}$

In QM. we talk about all (or but) distinct states are rep'd by a set of vectors:

$$|\psi_1\rangle \sim |\psi_2\rangle \quad \text{if} \quad |\psi_1\rangle = \lambda |\psi_2\rangle, \quad \lambda \in \mathbb{C}^*$$

$$\text{"rays"} \quad \underline{\psi} = \int e^{i\alpha} \psi, \quad \alpha \in \mathbb{R}$$

This is actually the projective complex

plane: $(z_1, z_2, z_3) \sim \lambda(z_1, z_2, z_3), \quad (\lambda \in \mathbb{C}^*) \in \mathbb{C}P^3$

represented by $[z_1 : z_2 : z_3]$

This is not a linear space:

$[1 : 0 : 0] \neq [0 : 1 : 0]$ cannot be defined.

no zero vector $[0 : 0 : 0]$

uniquely defined by the density matrix

$$P_\psi = \frac{|\psi\rangle\langle\psi|}{\langle\psi|\psi\rangle} \quad (P^2 = P)$$

projective Hilbert space $\mathbb{P}\mathcal{H} := (\mathcal{H} \setminus \{0\}) / \sim$

Consider symmetry operations on $\mathbb{P}\mathcal{H}$.

$$\text{The overlap. } O(P_1, P_2) = \text{Tr}(P_1 P_2) = \frac{|\langle\psi_1|\psi_2\rangle|^2}{\|\psi_1\|^2 \|\psi_2\|^2}$$

should be conserved.

But for all kinds of reasons we want to work on linear spaces \mathcal{H} . Aut(\mathcal{H})

Wigner's theorem, sym. operations are unitary or antiunitary

$$\begin{cases} \langle U\psi, U\phi \rangle = \langle \psi, \phi \rangle \\ \langle A\psi, A\phi \rangle = \overline{\langle \psi, \phi \rangle} = \langle \phi, \psi \rangle \end{cases}$$

$$(A^i = -iA, A = UK) \quad \hookrightarrow \text{plx. conj}$$

Now consider ^{unitary} symmetry operations on states.

one only needs

$$U(\beta_1) U(\beta_2) \psi = \underline{\alpha(\beta_1, \beta_2)} U(\beta_1 \beta_2) \psi$$

$$\alpha: G \times G \rightarrow U(1)$$

Not quite a group representation. projective rep.

$$\left(\begin{array}{l} \text{if } \alpha(g) = b(g) u(g) \\ \Rightarrow b(g) b(g_2) = f(g, g_2) b(g, g_2) \quad \forall g, g_2 \end{array} \right)$$

reduce back to rep.

rotation of spins:

$$\begin{array}{ccc} \text{classical} & & \text{quantum} \\ SO(3) & \longrightarrow & SU(2) \end{array}$$

SU(2) is a
spin projective
rep of SO(3)

Euler angles (xyz)

$$R(\phi, \theta, \psi) \rightarrow e^{i \frac{\phi}{2} \sigma^3} e^{i \frac{\theta}{2} \sigma^1} e^{i \frac{\psi}{2} \sigma^3}$$

$$R(2\pi, 0, 0) = 1 \quad \rightarrow \pm 1 \quad \text{for fermion / boson}$$

a \mathbb{Z}_2 phase

The central extension of G by A

$$1 \rightarrow A \xrightarrow{\iota} \tilde{E} \xrightarrow{\pi} G \rightarrow 1$$

is classified by the 2-cohomology

$$\text{group } H^2(G, A)$$

6. finite Heisenberg group.

$$P = \begin{pmatrix} 0 & 0 & 0 & 1 \\ 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \end{pmatrix} \quad Q = \begin{pmatrix} \omega & & & 0 \\ & \omega^2 & & \\ & & \omega^2 & \\ 0 & & & 1 \end{pmatrix}$$

$$\omega = e^{i \frac{2\pi}{d}}$$

$$\underline{QP = \omega PQ}$$

← Weyl relation of canonical commutation relation.

Some background.

$$[f, p] = i\hbar \quad (\hbar=1)$$

$$fP - Pf = i \quad P \cdot f \text{ acts on } f(P)$$

$$\Rightarrow A = e^{i\hbar p} \quad B = e^{i\eta f} \quad (\text{Weyl.})$$

$$AB = e^{i\hbar p} \cdot e^{i\eta f}$$

$$e^x e^y = e^z$$

$$z = x + y + \frac{1}{2}[x, y]$$

$$= e^{i(\hbar p + \eta f) + \frac{1}{2}[i\hbar p, i\eta f]}$$

$$+ \frac{1}{2}([x, [x, y]]$$

$$- [y, [x, y]])$$

$$+ \dots$$

$$BA = e^{i(\hbar p + \eta f) + \frac{1}{2}[i\eta f, i\hbar p]}$$

$$\Rightarrow AB = e^{i\hbar\eta} BA \quad (A, B = n \times n \text{ mats.})$$

$$\det(AB) = \omega^n \det(BA) \Rightarrow \underline{\omega^n = 1}$$

$$\begin{cases} A^k B = \omega^k B A^k \\ A B^l = \omega^l B^l A \end{cases}$$

$$\rightarrow A^k B^l = \omega^{kl} B^l A^k$$

$$k=n, d=1 \Rightarrow \underline{A^n} B = \omega^{\frac{1}{n}} B A^n \stackrel{?}{\Rightarrow} A^n = \mathbb{1}$$

$$\text{similarly } B^n = \mathbb{1}.$$

A general element in Heis_N has the form

$$\omega^a p^b q^c$$

$$(\omega^{a_1} p^{b_1} q^{c_1}) \cdot (\omega^{a_2} p^{b_2} q^{c_2}) = \omega^{a_3} p^{b_3} q^{c_3}$$

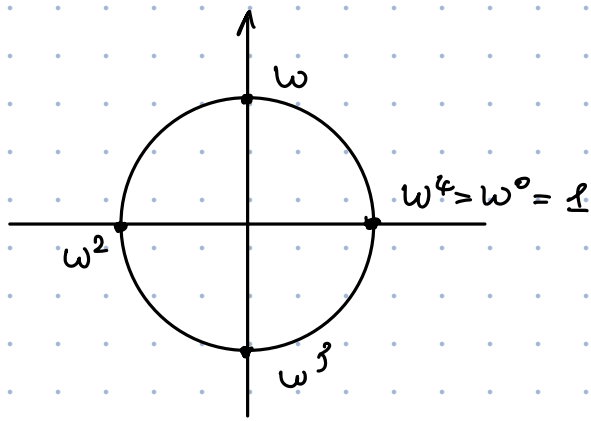
$$\begin{cases} a_3 = a_1 + a_2 + c_1 b_2 \\ b_3 = b_1 + b_2 \\ c_3 = c_1 + c_2 \end{cases}$$

$$\pi: \text{Heis}_N \longrightarrow \mathbb{Z}_N \times \mathbb{Z}_N$$

$$\omega^a p^b q^c \longmapsto (b \bmod N, c \bmod N)$$

$$\ker(\pi) = \left\{ \omega^a \cancel{p^{c_1}} \cancel{q^{c_2}} \right\} \cong \mathbb{Z}$$

$$\mathbb{1} \longrightarrow \mathbb{Z}_N \longrightarrow \text{Heis}_N \longrightarrow \mathbb{Z}_N \times \mathbb{Z}_N \longrightarrow \mathbb{1}$$



$$(P \cdot \Psi)(\omega^k) := \Psi(\omega^{k-1}) \quad \text{translation}$$

$$(Q \Psi)(\omega^k) := \omega^k \Psi(\omega^k) \quad \text{position operator}$$

$$(QP) \Psi(\omega^k) = \omega^k P \Psi(\omega^k) = \omega^k \Psi(\omega^{k-1})$$

$$(PQ) \Psi(\omega^k) = Q \Psi(\omega^{k-1}) = \omega^{k-1} \Psi(\omega^{k-1})$$

$$\Rightarrow QP = \omega PQ$$

$$N \rightarrow \infty \quad : \quad \mathbb{Z}_N \rightarrow U(1)$$

$$\mathbb{Z}_N \times \mathbb{Z}_N \rightarrow \mathbb{R} \times \mathbb{R}$$

$$1 \rightarrow U(1) \rightarrow \text{Heis}(\mathbb{R} \times \mathbb{R}) \rightarrow \mathbb{R} \times \mathbb{R} \rightarrow 1$$

7. More on group actions

7.1. Some defs and 3-0 theorem

Recall that the group action of G on a set X :

$$\phi : G \times X \rightarrow X$$

① left action. $\phi(g_1, \phi(g_2, x)) = \phi(g_1 g_2, x)$

$$g_1 \cdot (g_2 \cdot x) = (g_1 g_2) \cdot x$$

(right action $(x \cdot g_2) \cdot g_1 = x \cdot (g_2 g_1)$)

② $\phi(1_G, x) = x$

mention different forms of L & R actions
and induced actions on $\mathcal{F}[X \rightarrow Y]$

A G -action is:

see Moore's note

① effective: $\forall g \neq 1, \exists x \text{ s.t. } gx \neq x$

(ineffective $\exists g \neq 1, \forall x \text{ s.t. } gx = x$)

② transitive: $\forall x, y \in X, \exists g \text{ s.t. } \underline{y = g \cdot x}$

there is only one orbit

③ free: $\underline{\forall g \neq 1}, \underline{\forall x}, \underline{g \cdot x \neq x}$

Definitions.

1. isotropy group (stabilizer group)

$$\text{Stab}_G(x) := \{ g \in G \mid g \cdot x = x \} \subset G$$

$$(\equiv G^x)$$

$$\left(\begin{array}{l} g_1, g_2 \in G^x \\ g_1 \cdot x = x \quad g_2(g_1 \cdot x) = g_2 \cdot x = x \\ g_2 g_1 \in G^x \end{array} \right)$$

If the group action of G is free

$$\Leftrightarrow G^x = \{ 1 \} \quad \forall x \in X.$$

2. If $\exists g \in G^x \neq 1 \quad g \cdot x = x$. x is called a fixed point.

$$(X^g \equiv) \text{Fix}_X(g) = \{ x \in X \mid g \cdot x = x \} \subset X$$

is the fixed point set of g .

$$\text{free} \Leftrightarrow X^g = \emptyset \quad (g \neq 1)$$

3. $O_G(x) = \{ g \cdot x \mid \forall g \in G \}$

Theorem (Stabilizer-orbit)

Let X be a G -set. Each left-coset of

$G^x (\equiv \text{Stab}_G(x)) \quad (x \in X)$ is in a natural

1-1 correspondence with points in $D_G(x)$.

There exists a natural isomorphism

$$\begin{aligned}\varphi: D_G(x) &\longrightarrow G/G^x \\ g \cdot x &\longmapsto g \cdot G^x\end{aligned}$$

① Well defined.

$$g \cdot x = g' \cdot x$$

$$\Leftrightarrow (g'^{-1}g) \cdot x = x \Leftrightarrow g'^{-1}g \in G^x \Leftrightarrow g \cdot G^x = g' \cdot G^x$$

② surjective ✓

$$\text{injective: } g \cdot G^x = g' \cdot G^x \Rightarrow g \cdot x = g' \cdot x$$

$$\text{For a finite group: } |D_G(x)| = [G : G^x] = |G|/|G^x|$$

Example

1. G acts on G by conjugation $h \in G$.

$$D_G(h) = \{g h g^{-1}, \forall g \in G\} = C(h)$$

$$\text{Stab}_G(h) \equiv G = \{g \in G, \underline{g h g^{-1} = h}\} =: C_G(h)$$

Definition. The centralizer of h in G

$$C_G(h) := \{g \in G : g h = h g\}$$

(1) $C_G(h)$ is a subgroup

$$\textcircled{1} e \in C_G(h) : eh = he$$

$$\textcircled{2} \forall g_1, g_2 \in C_G(h) \quad (g_1 g_2^{-1})h = g_1 h g_2^{-1} = h g_1 g_2^{-1} \\ \Rightarrow (g_1 g_2^{-1}) \in C_G(h)$$

$$|C(h)| = [G : C_G(h)]$$

↑

number of conjugates of h

extend to subsets.

$$C_G(H) = \{ g \in G : gh = hg \quad \forall h \in H \}$$

$$C_G(G) = Z(G)$$

7.2. Practice with terminology of group actions.

1. $X = \{1, \dots, n\}$, $G = S_n$

① effective. \checkmark ($\forall \phi \neq 1, \exists x, \phi \cdot x \neq x$)

② transitive. \checkmark

③ free. \times ($\forall \phi \neq 1, \forall x, \phi \cdot x \neq x$)

keep j fixed. $\cong S_{n-1}$

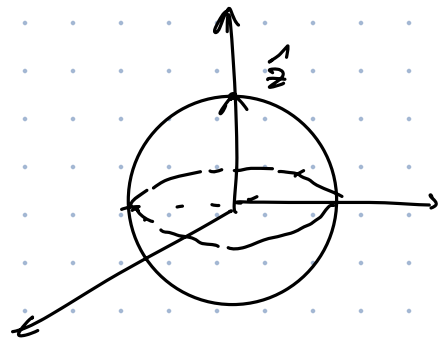
$$\cong S_{n-1}$$

2. $SO(3)$ acts on S^2

① effective. ✓

② transitive ✓

③ free? ✗



$$\text{Stab}_{SO(3)}\left(\frac{\hat{1}}{2}\right) = \left\{ \begin{pmatrix} \cos\phi & -\sin\phi & 0 \\ \sin\phi & \cos\phi & 0 \\ 0 & 0 & 1 \end{pmatrix}, \phi \in [0, 2\pi) \right\}$$

$$\cong SO(2)$$

$$\frac{\text{Orb}_{SO(3)}(\hat{n})}{\cong} \cong SO(3) / SO(2)_{\hat{n}}$$

$\cong S^2$